Synthesis, structure, and electronic structure of CsAgGa$_2$Se$_4$

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ABSTRACT

The new metal chalcogenide CsAgGa$_2$Se$_4$ has been synthesized by means of the reactive flux method. It crystallizes in the space group $P2_1/c$ of the monoclinic system with cell dimensions of $a=11.225(2)$ Å, $b=7.9443(16)$ Å, $c=21.303(4)$ Å, $\beta=103.10(3)$, $V=1850.3(6)$, and $Z=8$. The structure contains two-dimensional $^2_\infty[AgGa_2Se_4]^-$ layers separated by Cs$^+$ cations. The $^2_\infty[AgGa_2Se_4]^-$ superlayer possesses a novel chain–sublayer–chain structure: a $^2_\infty[Ag_2GaSe_6]^{7-}$ sublayer, composed of $^1_\infty[AgGaSe_4]^{4-}$ chains that are further connected by Ag$^+$ ions, is sandwiched by parallel $^1_\infty[Ga_3Se_8]^{7-}$ chains to generate the $^2_\infty[AgGa_2Se_4]^-$ superlayer. From a band structure calculation, the orbitals of all atoms have contributions to the bottoms of conduction bands, but the band gap is mainly determined by the 4s, 4p orbitals of Ga and Se.

1. Introduction

The ternary $M/M'/Q$ ($M=$coinage metal, $M'=$Al, Ga, In, $Q=$S, Se, Te) chalcogenides have been widely studied for their diverse structures and potential applications in nonlinear optics, solar energy conversion, etc. [1–5]. In contrast, research on the addition of an alkali metal to the ternary $M/M'/Q$ system has been very limited [6–8]. Introducing the strong electrostatic $A-Q$ ($A=$alkali metal) interactions in the covalently bonded $M-M'-Q$ framework could not only reduce the structure dimensionality from three-dimensional framework to two-dimensional layers or one-dimensional chains, but also alters the electronic structure and physical properties, such as the optical band gap.

So far, Na$_x$Cu$_{1-x}$In$_5$S$_8$ ($0<x<1$), NaAgGa$_6$Te$_{10}$, K$_2$CuIn$_3$Se$_6$, K$_2$CuGa$_3$Se$_6$, K$_2$AgIn$_3$Se$_6$ are the only compounds known in this quaternary $A/M/M'/Q$ ($A=$Na, K, Rb, Cs; $M=$coinage metal, $M'=$Al, Ga, In, $Q=$S, Se, Te) system [6–8]. Na$_x$Cu$_{1-x}$In$_5$S$_8$ ($0<x<1$) adopts the Al$_2$MgO$_4$ structure with the optical properties strongly dependent on the value of x. NaAgGa$_6$Te$_{10}$ adopts a three-dimensional structure composed of AgTe$_4$ and GaTe$_4$ tetrahedra with Na$^+$ in cavities. The isostructural K$_2$CuIn$_3$Se$_6$, K$_2$CuGa$_3$Se$_6$, and K$_2$AgIn$_3$Se$_6$ can be viewed as a stack of alternate anionic $[MM'_3Se_6]^-$ ($M=$Cu, Ag; $M'=$Ga, In) superlayers and K$^+$ cationic

layers. No compounds with the large alkali metal, namely Rb and Cs, were reported in this system. It has been shown that the size of alkali metal has a significant influence on the crystal structures [9–11]. Thus we carried out an exploration in the $A/M/M'/Q$ system, focusing on the larger alkali metals, and successfully discovered a new member in this family, CsAgGa$_2$Se$_4$, which contains $^2_\infty[AgGa_2Se_4]$ layers with a novel chain–sublayer–chain structure. In this paper, we report the synthesis, structure, and electronic structure of CsAgGa$_2$Se$_4$.

2. Experimental section

2.1. Single-crystal growth

In the first step, ternary AgGaSe$_2$ was prepared from high purity constituent elements by the high temperature solid-state synthesis technique. The starting materials were Ag (Sinopharm Chemical Reagent Co., Ltd., 99%), Ga (Sinopharm Chemical Reagent Co., Ltd., 99%), and Se (Sinopharm Chemical Reagent Co., Ltd., 99.9%), which were stored in an Ar-filled glovebox. The elements were loaded in the required stoichiometry into a silica ampule and sealed at a residual pressure of $<10^{-3}$ Pa and then reacted at 900 °C for 2 day.

Both binary metal polychalcogenides and metal halides could be used as reactive fluxes [12–16]. In this study, CsI was chosen for the reactive flux. The as-prepared AgGaSe$_2$ powder was mixed with the CsI (Sinopharm Chemical Reagent Co., Ltd., 99%) in the molar ratio of 1:2 and then loaded into a silica tube under an Ar

atmosphere in a glovebox, which was sealed under 10^{-3} Pa atmosphere and then placed in a computer-controlled furnace. The samples were heated to 800 °C in 20 h and kept at that temperature for 48 h, then cooled at a slow rate of 4 °C/h, and finally cooled to room temperature. The product consisted of yellow crystals of $CsAgGa_2Se_4$, which were manually selected for structure characterization. Analyses of the crystals with an EDX-equipped Hitachi S-3500 SEM showed the presence of Cs, Ag, Ga, and Se in the approximate molar ratio of 1:1:2:4. The crystals are stable in air.

2.2. Structure determination

Single-crystal X-ray diffraction data were collected with the use of graphite-monochromatized Mo $K\alpha$ ($\lambda = 0.71073$ Å) at 93 K on a Rigaku AFC10 diffractometer equipped with a Saturn CCD detector. Crystal decay was monitored by re-collecting 50 initial frames at the end of data collection and no detectable crystal decay was observed. The collection of the intensity data was carried out with the program Crystalclear [17]. Cell refinement and data reduction were carried out with the use of the program Crystalclear [17], and face-indexed absorption correction was performed numerically with the use of the program XPREP [18].

The structure was solved with Direct Methods implemented in the program SHELXS and refined with the least-squares program SHELXL of the SHELXTL.PC suite of programs [18]. The final refinement included anisotropic displacement parameters. The program STRUCTURE TIDY [19] was then employed to standardize the atomic coordinates. Additional details and structural data are given in Tables 1 and 2 and further information may be found in Supplementary material.

2.3. Band structure calculation

The electronic properties are calculated using the plane-wave pseudopotential method [20] implemented in the CASTEP package [21]. The local density functional (LDA) with a high kinetic-energy cutoff of 650 eV is adopted for all the calculations. The preconditioned conjugated gradient (CG) band-by-band method [22] used in CASTEP ensures a robust efficient search of the energy minimum of the electronic structure ground state. The optimized normal-conserving pseudopotentials [23] in Kleinman–Bylander form [24] for Ag, Cs, Ga, and Se allow us to use a small plane-wave

Table 1
Crystal data and structure refinement for $CsAgGa_2Se_4$.

	$CsAgGa_2Se_4$
fw	696.06
a (Å)	11.225(2)
b (Å)	7.9443(16)
c (Å)	21.303(4)
β (deg.)	103.10(3)
Space group	$P2_1/c$
V (Å^3)	1850.3(6)
Z	8
T (K)	93 (2)
λ (Å)	0.71073
ρ_c (g/cm^3)	4.998
μ (cm^{-1})	274.04
$R(F)^a$	0.0479
$R_w (F_o^2)^b$	0.1188

^{a} $R(F) = \sum||F_o| - |F_c||/\sum|F_o|$ for $F_o^2 > 2\sigma(F_o^2)$.
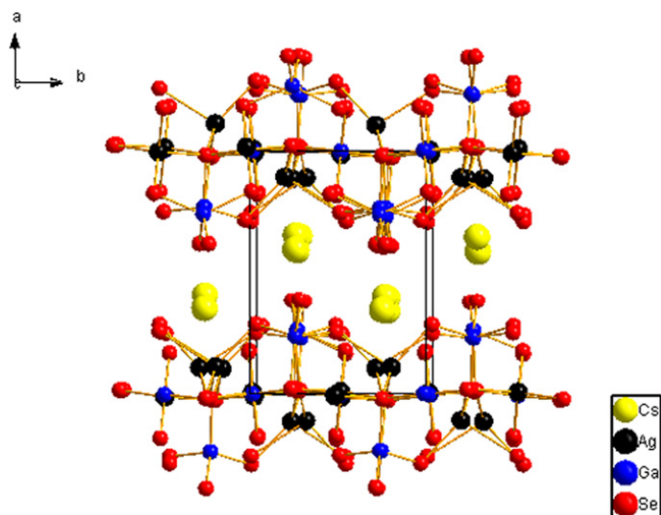^{b} $R_w(F_o^2) = \{\sum[w(F_o^2 - F_c^2)^2]/\sum wF_o^4\}^{1/2}$ for all data.

Table 2
Selected bond lengths (Å) for $CsAgGa_2Se_4$.

Atoms	Distances	Atoms	Distances	Atoms	Distances
Ag_1–Se_4	2.577(2)	Ga_1–Se_8	2.409(3)	Ga_3–Se_1	2.392(3)
Ag_1–Se_8	2.590(2)	Ga_1–Se_3	2.409(3)	Ga_3–Se_8	2.419(3)
Ag_1–Se_2	2.663(3)	Ga_1–Se_2	2.413(3)	Ga_3–Se_6	2.438(3)
Ag_1–Se_3	2.666(3)	Ga_1–Se_4	2.415(3)	Ga_3–Se_5	2.462(3)
Ag_2–Se_8	2.607(3)	Ga_2–Se_7	2.388(3)	Ga_4–Se_3	2.400(3)
Ag_2–Se_4	2.629(2)	Ga_2–Se_6	2.424(3)	Ga_4–Se_2	2.411(3)
Ag_2–Se_5	2.633(3)	Ga_2–Se_4	2.426(3)	Ga_4–Se_7	2.435(3)
Ag_2–Se_6	2.972(3)	Ga_2–Se_5	2.462(3)	Ga_4–Se_1	2.446(3)

basis set without compromising the accuracy required by our study. The electrons below 4s and 5s orbitals for sliver and cesium are treated as the core electrons. For gallium the 3d, 4s and 4p electrons are chosen as the valence electrons and for selenium, they are 4s and 4p electrons. Monkhorst–Pack k-point meshes [25] with a density of $(1 \times 2 \times 1)$ points in the Brillouin zone of the $CsAgGa_2Se_4$ unit cell are chosen.

3. Results and discussion

3.1. Structure

$CsAgGa_2Se_4$ crystallizes in a new structure type in space group $P2_1/c$ of the monoclinic system. The asymmetric unit contains two crystallographically independent Cs atoms, two independent Ag atoms, four independent Ga atoms, and eight independent Se atoms. The coordination polyhedra of Cs1 and Cs2 cations are irregular with coordination numbers 8 and 9 and Cs–Se distances in the range of 3.573(2)−4.260(3) Å. These Cs–Se distances are close to those in $Cs_{10}P_8Se_{20}$ [26] (3.421(6)–4.240(2) Å), $CsAuUSe_3$ [27] (3.604(3)–3.719(3) Å), and $Cs_2Hg_2USe_5$ [28] (3.536(2)–4.224(2) Å). Each Ag atom is coordinated to a slightly distorted tetrahedron of four Se atoms. The Ag–Se distances range from 2.577(2) to 2.972(3) Å, which are comparable to those of $Ba_2Cu_{1.78(7)}Ag_{2.22}Se_5$ [29] (2.552(2)–3.096(2) Å), $BaAg_2SnSe_4$ [30] (2.5670(5)–2.9334(6) Å), $CsTb_2Ag_3Se_5$ [31] (2.6135(8)–2.766(1) Å), and $AgBiP_2Se_6$ [32] (2.762(2)–3.032(3) Å). Each Ga atom is also coordinated to a slightly distorted tetrahedron of four Se atoms. The Ga–Se distances of 2.388(3)–2.462(3) Å are in good agreement with those in $BaGa_4Se_7$ [33] (2.361(2)–2.488(2) Å). Since there are no Se–Se bonds in the structures, the oxidation states of 1+, 1+, 3+, and 2− can be assigned to Cs, Ag, Ga, and Se, respectively.

The crystal structure of the compound can be considered as a stack of alternate anionic $^2_\infty[AgGa_2Se_4]$ superlayers and Cs^+ cationic layers (Fig. 1). The structural novelty of $CsAgGa_2Se_4$ derives from the fact that its anionic superlayer is composed of a sublayer sandwiched by parallel chains, as shown in Fig. 2. The two-dimensional $^2_\infty[Ag_2GaSe_6]_{7-}$ sublayer is assembled from $^1_\infty[AgGaSe_4]_{4-}$ chains that are connected by Ag atoms (Fig. 3(a)) and the parallel $^1_\infty[Ga_3Se_8]_{7-}$ chains running above or below the sublayer is composed of $GaSe_4$ tetrahedra (Fig. 3(b)). The structure of $CsAgGa_2Se_4$ is different from that of $K_2AgIn_3Se_6$, which also features a chain–sublayer–chain structure, in two aspects. Firstly, the $In1Se_4$, $In2Se_4$, and $AgSe_4$ are connected to each other by corner-sharing only in the sublayer of $K_2AgIn_3Se_6$ (Fig. 4(a)), while the sublayer in $CsAgGa_2Se_4$ is more complicated in that the edge-sharing $Ag1Se_4$ and $Ga1Se_4$ tetrahedra form $^1_\infty[AgGaSe_4]_{4-}$ chains along the b direction, which are further connected by the $Ag2Se_4$ tetrahedra by corner-sharing to form the $^2_\infty[Ag_2GaSe_6]_{7-}$ sublayer (Fig. 3(a)); Secondly, in the $K_2AgIn_3Se_6$ structure, the chain above or below the sublayer is formed by the corner-sharing $InSe_4$ tetrahedra, as shown in Fig. 4(b), while in the

Fig. 1. Unit cell of the CsAgGa$_2$Se$_4$ structure.

Fig. 2. $^2_\infty$[AgGa$_2$Se$_4$]$^-$ superlayer of CsAgGa$_2$Se$_4$ structure.

CsAgGa$_2$Se$_4$ structure, those chains are formed by both corner-sharing and edge-sharing GaSe$_4$ tetrahedra (Fig. 3(b)).

3.2. Band structure calculation

The calculated band structures of the CsAgGa$_2$Se$_4$ in the unit cell are plotted along the symmetry lines in Fig. 5. It is shown that CsAgGa$_2$Se$_4$ is a direct gap crystal, and the band gap at *G* is 1.84 eV, which is larger than the calculated band gaps of AgGaSe$_2$ (1.42 eV) [34] using the same method. It is known that the band gap calculated by LDA is usually smaller than the experimental data due to the discontinuity of exchange-correlation energy. Typically, the difference between the experimental and calculated band gaps is ranged from 0.2 to 0.5 eV for semiconductor systems. Thus, we estimate that the experimental band gap of CsAgGa$_2$Se$_4$ is about 2.0–2.3 eV, which is consistent with the yellow color of the crystal. In comparison, the AgGaSe$_2$ crystal is red in color, which also indicate that AgGaSe$_2$ possess a smaller band gap than CsAgGa$_2$Se$_4$.

Fig. 6 gives the total density of states (DOS) and partial DOS (PDOS) projected on the constitutional atoms of the CsAgGa$_2$Se$_4$ crystal, in which several characteristics can be deduced. (1) The bands localized at −85 eV (not shown) and −14 eV mostly consist by 4*s*, 4*p* orbitals of Ag, 5*s* orbitals of Cs, and 3*d* orbitals of Ga. They do not form the chemical bonding with other orbitals. (2) The orbitals of Cs contribute nothing to the upper region of the valence bands (VB) from −10 eV to 0 eV. These electronic bands
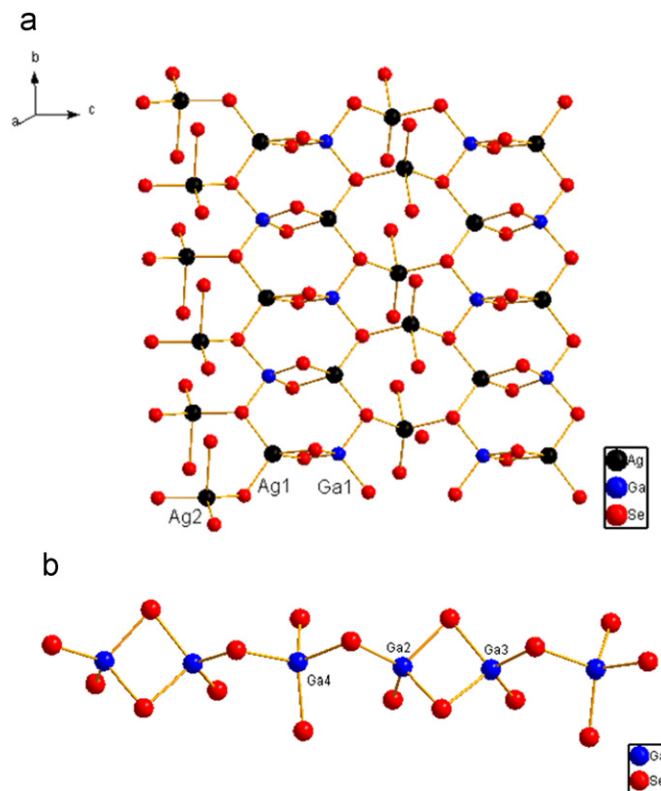

Fig. 3. (a) Sublayer of the CsAgGa$_2$Se$_4$ structure and (b) the chain above or below the sublayer of the CsAgGa$_2$Se$_4$ structure
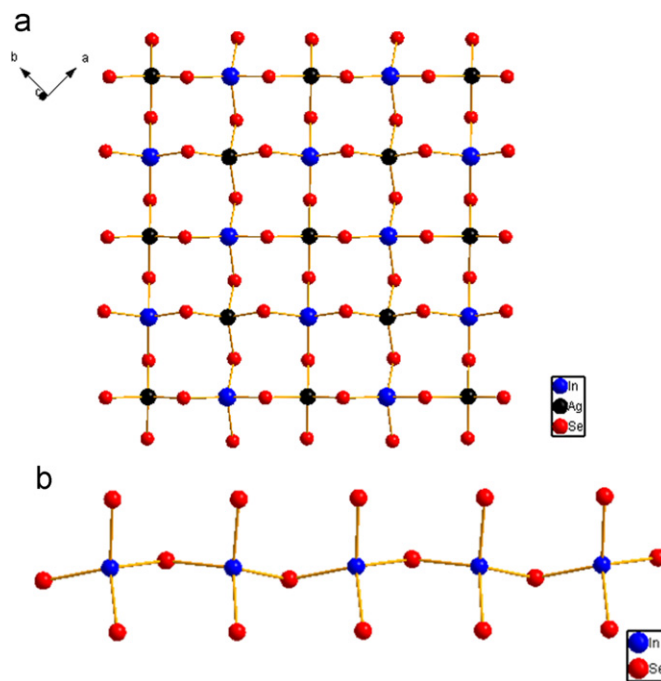

Fig. 4. (a) Sublayer of the K$_2$AgIn$_3$Se$_6$ structure and (b) the chain above or below the sublayer of the K$_2$AgIn$_3$Se$_6$ structure.

are mainly composed of Ag 4*d*, Ga 4*s* and 4*p* orbitals, as well as Se 4*p* orbitals. The quite large mixture among their states is clearly displayed, revealing that the *Ga–Se*, *Ag–Se* bonds are strong. The very top of the VB from −4 to 0 eV mostly consists of the Se 4*p* orbitals. (3) The orbitals of all atoms have contributions to the bottoms of conduction bands, but compared with AgGaSe$_2$, whose
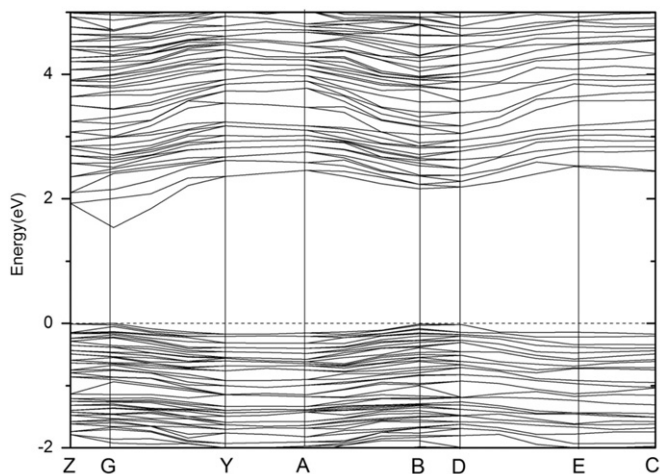
Fig. 5. Band structure of CsAgGa$_2$Se$_4$ along the lines of high symmetry points in the Brillouin zone. The dash line indicates the VB maximum.
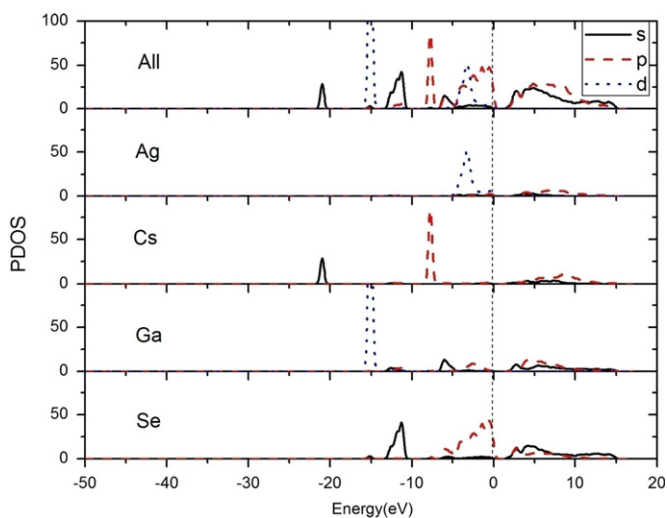

Fig. 6. Partial density of states of CsAgGa$_2$Se$_4$. The solid, dot-dash, and dash lines are the *s*, *p*, and *d* orbitals, respectively. The broken vertical lines indicate the VB maximum.

band gap is mainly determined by the orbitals of Ag and Se [34], the addition of Cs made the orbital of Ga contribute more than Ag to the band gap. The band gap of CsAgGa$_2$Se$_4$ is mainly determined by the 4*s*, 4*p* orbitals of Ga and Se. This may be the reason for the increased band gap of CsAgGa$_2$Se$_4$ compared with AgGaSe$_2$.

4. Conclusions

The new selenide CsAgGa$_2$Se$_4$ has been synthesized for the first time from the reaction between AgGaSe$_2$ and CsI. It crystallizes in a new structure type in the monoclinic space group $P2_1/c$. The structure contains two-dimensional $^2_\infty$[AgGa$_2$Se$_4$]$^-$ layers separated by Cs$^+$ cations. The novel $^2_\infty$[AgGa$_2$Se$_4$]$^-$ layer consists of a $^2_\infty$[Ag$_2$GaSe$_6$]$^{7-}$ sublayer sandwiched by $^1_\infty$[Ga$_3$Se$_8$]$^{7-}$ chains. The $^2_\infty$[Ag$_2$GaSe$_6$]$^{7-}$ sublayers is formed by corner and edge-sharing AgSe$_4$ and GaSe$_4$ tetrahedra and the parallel $^1_\infty$[Ga$_3$Se$_8$]$^{7-}$ chains above or below the sublayer is composed of corner and edge-sharing GaSe$_4$ tetrahedra. The electronic structure calculation indicates that, compared with AgGaSe$_2$, whose band gap is

mainly determined by the orbitals of Ag and Se, the orbital of Ga contribute more than Ag to the optical band gap in CsAgGa$_2$Se$_4$. Thus the introducing of Cs not only breaks the three-dimensional framework of AgGaSe$_2$ into a two-dimensional structure but also alters the electronic structure.

Appendix A. Supplementary materials

Supplementary data associated with this article can be found in the online version at doi:10.1016/j.jssc.2011.11.014.

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